## OSCILLATIONS

## Periodic Motion and Oscillatory motion

If a body repeats its motion along a certain path, about a fixed point, at a definite interval of time, it is said to have a periodic motion
If a body moves to and fro, back and forth, or up and down about a fixed point in a definite interval of time, such motion is called an oscillatory motion. The body performing such motion is called an oscillator.

## Simple Harmonic Motion

The periodic motion of a body about a fixed point, on a linear path, under the influence of the force acting towards the fixed point and proportional to displacement of the body from the fixed point is called a simple harmonic motion (SHM) A body performing simple harmonic motion is known as Simple Harmonic Oscillator (SHO) Simple harmonic motion is a special type of periodic motion in which
(i) The particle oscillates on a straight line.
(ii) The acceleration of the particle is always directed towards a fixed point on the straight line.
(iii) The magnitude of acceleration is proportional to the displacement of the particle from fixed point.


This fixed point is called the centre of the oscillation or mean position. Taking this point as origin " $O$ ".
The maximum displacement of oscillator on either side of the mean position is called amplitude denoted by A
The time required to complete one oscillation is known as periodic time $(T)$ of oscillator In other words, the least time interval of time after which the periodic motion of an oscillator repeat itself is called a periodic time of the oscillator. Distance travelled by oscillator is 4A in periodic time
The number of oscillations completed by simple harmonic oscillator in one second is defined as frequency. SI unit is $\mathrm{s}^{-1}$ or hertz (H)
It is denoted by $f$ and $f=1 / T$
$2 \pi$ times the frequency of an oscillator is called the angular frequency of the oscillator It is denoted by $\omega$. Its SI unit is rad $\mathrm{s}^{-1}, \omega=2 \pi f$.
If we draw the graph of displacement of SHO against time as shown in figure, which is as shown in figure

## PHYSICS NOTES



Mathematical equation is the function of time
$x(t)=A \sin (\omega t+\phi)----e q(1)$
Here $X(t)$ represents displacement at time $t$
A = Amplitude
$(\omega t+\phi)=$ Phase
$\phi=$ Initial phase (epoch) for a given graph $\phi=0$.
If oscillation have started from negative $x$ end (M) then initial phase is $-\pi / 2$
If oscillation stated from positive end ( N ) then initial phase is $\pi / 2$
$\omega$ = angular frequency

## Velocity

Velocity of oscillator is

$$
v(t)=\frac{d x(t)}{d t}
$$

$v(t)=\omega A \cos (\omega t+\varphi)$

$$
\begin{gathered}
v(t)= \pm \omega A\left(\sqrt{1-\sin ^{2}(\omega \mathrm{t}+\varphi)}\right) \\
v(t)= \pm \omega\left(\sqrt{A^{2}-A^{2} \sin ^{2}(\omega \mathrm{t}+\varphi)}\right)
\end{gathered}
$$

From eq(1)

$$
v(t)= \pm \omega\left(\sqrt{A^{2}-x^{2}}\right)
$$

Velocity is maximum at $x=0$ or equilibrium position $v= \pm \omega \mathrm{A}$
Velocity is minimum at $x=A$ or extreme positions $v=0$
Note that velocity is out of phase of displacement by $\pi / 2$

## Acceleration

$$
a(t)=\frac{d v(t)}{d t}
$$

$a(t)=-\omega^{2} A \sin (\omega t+\phi)$
$a(\mathrm{t})=-\omega^{2} \mathrm{x}(\mathrm{t})$
At $\mathrm{x}=0$, acceleration is $\mathrm{a}=0$

## PHYSICS NOTES

And maximum at $x= \pm A, a=\mp \omega^{2} A$
Note velocity is out of phase of displacement by $\pi$

## Solved Numerical

Q) A particle moving with S.H.M in straight line has a speed of $6 \mathrm{~m} / \mathrm{s}$ when 4 m from the centre of oscillations and a speed of $8 \mathrm{~m} / \mathrm{s}$ when 3 m from the centre. Find the amplitude of oscillation and the shortest time taken by the particle in moving from the extreme position to a point midway between the extreme position and the centre

## Solution

From the formula for velocity

$$
\begin{gathered}
v(t)= \pm \omega\left(\sqrt{A^{2}-x^{2}}\right) \\
6= \pm \omega\left(\sqrt{A^{2}-4^{2}}\right)
\end{gathered}
$$

And

$$
8= \pm \omega\left(\sqrt{A^{2}-3^{2}}\right)
$$

By taking ratio of above equations

$$
\frac{6}{8}=\frac{\sqrt{A^{2}-4^{2}}}{\sqrt{A^{2}-3^{2}}}
$$

On simplifying we get $A=5 \mathrm{~m}$
On substituting value of $A$ in equation

$$
\begin{aligned}
& 6= \pm \omega\left(\sqrt{A^{2}-4^{2}}\right) \\
& 6= \pm \omega\left(\sqrt{5^{2}-4^{2}}\right)
\end{aligned}
$$

$\omega=2 \mathrm{rad}$
From the equation for oscillation
$x(t)=A \sin (\omega t+\phi)$

$$
\begin{gathered}
\frac{A}{2}=A \sin 2 t \\
0.5=\sin 2 t \\
\Rightarrow 2 t=\pi / 6
\end{gathered}
$$


$t=\pi / 12$ this is the time taken by oscillator to move from centre to midway
Now Time taken by oscillator to move from centre to extreme position is T/4
As $\mathrm{T}=2 \pi / \omega$ thus
Time taken to reach to extreme position from the centre $=$

$$
\frac{T}{4}=\frac{2 \pi}{4 \omega}=\frac{2 \pi}{4 \times 2}=\frac{\pi}{4}
$$

## PHYSICS NOTES

Now time taken to reach oscillator to reach middle from extreme point is $=$

$$
\frac{\pi}{4}-\frac{\pi}{12}=\frac{\pi}{6} \sec
$$

Q) A point moving in a straight line SHM has velocities $\mathrm{v}_{1}$ and $\mathrm{v}_{2}$ when its displacements from the mean position are $x_{1}$ and $x_{2}$ respectively. Show that the time period is

$$
2 \pi \sqrt{\frac{x_{1}^{2}-x_{2}^{2}}{v_{2}^{2}-v_{1}^{2}}}
$$

## Solution

From the formula for velocity

$$
\begin{gathered}
v(t)= \pm \omega\left(\sqrt{A^{2}-x^{2}}\right) \\
v_{1}^{2}=\omega^{2}\left(A^{2}-x_{1}^{2}\right) \\
A^{2}=\frac{v_{1}^{2}}{\omega^{2}}+x_{1}^{2}---e q(1) \\
A^{2}=\frac{v_{2}^{2}}{\omega^{2}}+x_{2}^{2}---e q(2)
\end{gathered}
$$

From equation (1) and (2)

$$
\begin{gathered}
\frac{v_{1}^{2}}{\omega^{2}}+x_{1}^{2}=\frac{v_{2}^{2}}{\omega^{2}}+x_{2}^{2} \\
\frac{v_{2}^{2}}{\omega^{2}}-\frac{v_{1}^{2}}{\omega^{2}}=x_{1}^{2}-x_{2}^{2} \\
\frac{1}{\omega^{2}}\left(v_{2}^{2}-v_{1}^{2}\right)=x_{1}^{2}-x_{2}^{2} \\
\frac{T^{2}}{(2 \pi)^{2}}\left(v_{2}^{2}-v_{1}^{2}\right)=x_{1}^{2}-x_{2}^{2} \\
T=2 \pi \sqrt{\frac{x_{1}^{2}-x_{2}^{2}}{v_{2}^{2}-v_{1}^{2}}}
\end{gathered}
$$

Relation between simple harmonic motion and uniform circular motion



$X=A \cos (\omega t+\phi)$
i.e. P moves with simple harmonic motion

Thus, when a particle moves with uniform circular motion, its projection on a diameter moves with simple harmonic motion. The angular frequency $\omega$ of simple harmonic motion is the same as the angular speed of the reference point.
The velocity of $Q$ is $v=\omega A$. The component of $v$ along the $x$-axis is
$V_{x}=-v \sin (\omega t+\phi)$
$V_{x}=-\omega A \sin (\omega t+\phi)$,
Which is also the velocity of $p$. The acceleration of $Q$ is centripetal and has a magnitude, $\mathrm{a}=\omega^{2} \mathrm{~A}$
The component of ' $a$ ' along the $x$-axis is
$A_{x}=-\operatorname{acos}(\omega t+\phi)$,
$A_{x}=-\omega^{2} A \cos (\omega t+\phi)$,
Which is the acceleration of $P$

## The force law for simple harmonic motion

We know that
$\mathrm{F}=\mathrm{ma}$
As $a=-\omega^{2} x(t)$
$\mathrm{F}=-\mathrm{m} \omega^{2} \mathrm{x}(\mathrm{t})$
This force is restoring force
According to Hook's law, the restoring force is given by
$\mathrm{F}=-\mathrm{kx}(\mathrm{t})$
With $k$ as spring constant
Thusk $=\mathrm{m} \omega^{2}$
In figure shown $Q$ is a point moving on a circle of radius A with constant angular speed $\omega$ (in rad/sec). P is the perpendicular projection of $Q$ on the horizontal diameter, along the $x$-axis.
Let us take $Q$ as the reference point and the circle on which it moves the reference circle. As the reference point revolves, the projected point $P$ moves back and forth along the horizontal diameter

Let the angle between the radius $O Q$ and the $x$-axis at the time $t=0$ be called $\phi$. At any time $t$ later the angle between $O Q$ and the $x$-axis is ( $\omega t+\phi$ ), the point $Q$ moving with constant angular speed $\omega$. The $x$ coordinate of $Q$ at any time is, therefore
$\therefore$ angular frequency

$$
\omega=\sqrt{\frac{k}{m}}
$$

And frequency of oscillation

$$
f=\frac{1}{2 \pi} \sqrt{\frac{k}{m}}
$$

In many cases, the simple harmonic motion can also occur even without spring. In that case $k$ is called the force constant of SHM and it is restoring force per unit displacement ( $\mathrm{K}=-\mathrm{F} / \mathrm{x}$ )

## Solved numerical

Q) A spring balance has a scale that reads 50 kg . The length of the scale is 20 cm . A body suspended from this spring, when displaced and released, oscillates with period of 0.6 s . Find the weight of the body
Solution:
Here $m=50 \mathrm{~kg}$
Maximum extension of spring $x=20-0=20 \mathrm{~cm}=0.2 \mathrm{~m}$
Periodic time $\mathrm{T}=0.6 \mathrm{~s}$
Maximum force $F=m g$
$\mathrm{F}=50 \times 9.8=490 \mathrm{~N}$
$K=F / X=490 / 0.2=2450 \mathrm{~N} \mathrm{~m}^{-1}$
As

$$
\begin{gathered}
T=2 \pi \sqrt{\frac{m}{k}} \\
m=\frac{(0.6)^{2} \times 2450}{4 \times(3.14)^{2}}=22.36 \mathrm{~kg}
\end{gathered}
$$

$\therefore$ Weight of the body $=\mathrm{mg}=22.36 \times 9.8=219.1 \mathrm{~N}=22.36 \mathrm{kgf}$

## Examples of simple harmonic motion

## Simple pendulum

A simple pendulum consists of a heavy particle suspended from a fixed support through a light, inextensible and torsion less string
The time period of simple pendulum can be found by force or torque method and also by energy method
(a)Force method: the mean position or the equilibrium position of the simple pendulum is when $\theta=0$ as shown in figure(i). the length of the string is $I$, and mass of the bob is $m$ When the bob is displaced through distance ' $x$ ', the forces acting on it are shown in figure(ii)


The restoring force acting on the bob to bring it to the mean position is
$\mathrm{F}=-\mathrm{mg} \sin \theta$ (-ve sign indicates that force is directed towards the mean position)
For small angular displacement
$\sin \theta \approx \theta=x / l$

$$
\begin{aligned}
& \therefore F=-m g \frac{x}{l} \\
& \therefore a=-g \frac{x}{l}
\end{aligned}
$$

Comparing it with equation of simple harmonic motion $a=-\omega^{2} x$

$$
\begin{aligned}
\omega^{2} & =\frac{g}{l} \\
\therefore \omega & =\sqrt{\frac{g}{l}}
\end{aligned}
$$

Time period

$$
T=\frac{2 \pi}{\omega}=2 \pi \sqrt{\frac{l}{g}}
$$

## (b)Torque method:

Now taking moment of force acting on the bob about 0

$$
\tau=-(m g \sin \theta) l---e q(1)
$$

Also from Newton's second law

$$
\tau=I \alpha---e q(2)
$$

From equation(1) and (2)

$$
I \alpha=-(m g \sin \theta) l
$$

Since $\theta$ is small

$$
I \alpha=-(m g \theta) l
$$

But Moment of inertia $\mathrm{I}=\mathrm{ml}^{2}$

$$
\begin{gathered}
m l^{2} \alpha=-(m g \theta) l \\
\therefore \alpha=-\left(\frac{g}{l}\right) \theta
\end{gathered}
$$

Comparing with simple harmonic motion equation , $\alpha=-\omega^{2} \theta$

$$
\begin{gathered}
\omega=\sqrt{\frac{g}{l}} \\
T=\frac{2 \pi}{\omega}=2 \pi \sqrt{\frac{l}{g}}
\end{gathered}
$$

Note: component of force $m g \cos \theta$ cannot produce torque because it passes through fixed point

## (c) Energy method:

Let the potential energy at the mean position be zero. Let the bob is displaced through an angle ' $\theta$ '. Let its velocity be ' $v$ '


Then potential energy at the new position
$\mathrm{U}=\mathrm{mgl}(1-\cos \theta)$
Kinetic energy at this instant $K=(1 / 2) \mathrm{mv}^{2}$
Total mechanical energy at this instant
$E=U+K$
$E=m g l(1-\cos \theta)+(1 / 2) m v^{2}$
We know, in simple harmonic motion $\mathrm{E}=$ constant

$$
\begin{array}{r}
\frac{d E}{d t}=0 \\
\Rightarrow m g l\left[\sin \theta \frac{d \theta}{d t}\right]+m v \frac{d v}{d t}=0
\end{array}
$$

$$
\begin{gathered}
\text { Bur } v=\omega l \\
\therefore v=l \frac{d \theta}{d t} \\
\therefore m g l\left[\sin \theta \frac{d \theta}{d t}\right]+m l \frac{d \theta}{d t} \frac{d v}{d t}=0 \\
\therefore g[\sin \theta]+\frac{d v}{d t}=0
\end{gathered}
$$

$$
\frac{d v}{d t}=a=-g[\sin \theta] \approx g \theta
$$

But $\theta=x / I$

$$
\therefore a=-g \frac{x}{l}
$$

But $a=-\omega^{2} x$

$$
\begin{aligned}
& \therefore \omega^{2}=\frac{g}{l} \\
& \omega=\sqrt{\frac{g}{l}}
\end{aligned}
$$

Periodic time

$$
T=\frac{2 \pi}{\omega}=2 \pi \sqrt{\frac{l}{g}}
$$

## (a)Simple pendulum in a lift:

If the simple pendulum is oscillating in a lift moving with acceleration a, then the effective $g$ of the pendulum is
$g_{\text {eff }}=g \pm a$

+ sign is taken when lift is moving upward
- ve sign is taken when lift is moving downward

Hence periodic time of pendulum

$$
T=2 \pi \sqrt{\frac{l}{g \pm a}}
$$

## (b)Simple pendulum in the compartment of a train

If the simple pendulum is oscillating in a compartment of a train accelerating or retarding horizontally at the rate ' $a$ ' then the effective value of $g$ is

$$
g_{e f f}=\sqrt{g^{2}+a^{2}}
$$

Hence periodic time of pendulum

$$
T=2 \pi \sqrt{\frac{l}{\sqrt{g^{2}+a^{2}}}}
$$

## (c)Seconds pendulum

The pendulum having the time-period of two seconds, is called the second pendulum. It takes one second to go from one end to the other end during oscillation. It also crosses the mean position at every one second

## Solved Numerical

Q) Length of a second's pendulum on the surface of earth is $I_{1}$ and $I_{2}$ at height ' $h$ ' from the surface of earth. Prove that the radius of the earth is given by

$$
R_{e}=\frac{h \sqrt{l_{2}}}{\sqrt{l_{1}}-\sqrt{l_{2}}}
$$

Solution:
Period of second's pendulum is 2 sec

$$
2=2 \pi \sqrt{\frac{l_{1}}{g}}
$$

And $g$ at surface of earth is given by

$$
\begin{gathered}
g=\frac{G M_{e}}{R_{e}^{2}} \\
\therefore 2=2 \pi \sqrt{\frac{l_{1} R_{e}^{2}}{G M_{e}}}---e q(1)
\end{gathered}
$$

$g$ at height $h$ from surface of earth is

$$
g=\frac{G M_{e}}{\left(R_{e}+h\right)^{2}}
$$

Thus at height h

$$
2=2 \pi \sqrt{\frac{l_{2}\left(R_{e}+h\right)^{2}}{G M_{e}}}---e q(2)
$$

From equation(1) and (2) we get

$$
\begin{gathered}
2 \pi \sqrt{\frac{l_{2}\left(R_{e}+h\right)^{2}}{G M_{e}}}=2 \pi \sqrt{\frac{l_{1} R_{e}^{2}}{G M_{e}}} \\
l_{2}\left(R_{e}+h\right)^{2}=l_{1} R_{e}^{2} \\
\sqrt{l_{2}}\left(R_{e}+h\right)=\sqrt{l_{1}} R_{e} \\
\sqrt{l_{2}} h=R_{e}\left(\sqrt{l_{1}}-\sqrt{l_{2}}\right) \\
R_{e}=\frac{h \sqrt{l_{2}}}{\sqrt{l_{1}}-\sqrt{l_{2}}}
\end{gathered}
$$

## Combinations of springs

## Series combination



As show in figure consider a series combination of two massless spring of spring constant $k_{1}$ and $k_{2}$

In this system when the combination of two springs is displaced to a distance $y$, it produces extension $x_{1}$ and $x_{2}$ in two springs of force constants $k_{1}$ and $k_{2}$.
$F=-k_{1} x_{1} ; F=-k_{2} x_{2}$
where $F$ is the restoring force.

$$
x=x_{1}+x_{2}=-F\left[\frac{1}{k_{1}}+\frac{1}{k_{2}}\right]
$$

We know that $F=-k x$
$x=-F / k$
From the above equations,

$$
\begin{aligned}
& \frac{1}{k}=\frac{1}{k_{1}}+\frac{1}{k_{2}} \\
& k=\frac{k_{1} k_{2}}{k_{1}+k_{2}}
\end{aligned}
$$

Time period $T$

$$
T=2 \pi \sqrt{\frac{m\left(k_{1}+k_{2}\right)}{k_{1} k_{2}}}
$$

Frequency f

$$
f=\frac{1}{2 \pi} \sqrt{\frac{k_{1} k_{2}}{m\left(k_{1}+k_{2}\right)}}
$$

If both the springs have the same spring constant, $k_{1}=k_{2}=k$. then equivalent spring constant $k^{\prime}=k / 2$

$$
\begin{gathered}
T=2 \pi \sqrt{\frac{m(2)}{k}} \\
f=\frac{1}{2 \pi} \sqrt{\frac{k^{\prime}}{m}}=\frac{1}{2 \pi} \sqrt{\frac{k}{2 m}}
\end{gathered}
$$

## Parallel combination

(i) Consider a situation as shown in figure where a body of mass $m$ is attached in between the two massless springs of spring constant $\mathrm{k}_{1}$ and $\mathrm{k}_{2}$. let the body is left free for SHM in
 vertical plane after pulling mass $m$ In this situation when a body is pulled lower through small displacement $y$. lower spring gets compressed by $y$, while upper spring elongate by $y$. hence restoring forces $F_{1}$ and $F_{2}$ set up in both these springs will act in the same direction.
Net restoring force will be
$\mathrm{F}=\mathrm{F}_{1}+\mathrm{F}_{2}$
$F=-k_{1} y-k_{2} y$
$F=-\left(k_{1}+k_{2}\right) y$
If $\mathrm{k}^{\prime}$ is the is equivalent spring constant then
$F=-k^{\prime} y$ thus
$k^{\prime}=k_{1}+k_{2}$
Now periodic time T

$$
T=2 \pi \sqrt{\frac{m}{k^{\prime}}}=\sqrt{\frac{m}{k_{1}+k_{2}}}
$$

Frequency

$$
f=\frac{1}{2 \pi} \sqrt{\frac{k^{\prime}}{m}}=\frac{1}{2 \pi} \sqrt{\frac{k_{1}+k_{2}}{m}}
$$

If $\mathrm{k}_{1}=\mathrm{k}_{2}=\mathrm{k}$ then

$$
T=2 \pi \sqrt{\frac{m}{2 k}}
$$

(ii) Two massless springs of equal lengths having force constants $k_{1}$ and $k_{2}$ respectively are suspended vertically from a rigid support as shown in figure. At their free ends, a block of
 mass $m$ having non-uniform density distribution is suspended so that spring undergoes equal extension In this situation two bodies are pulled down through a small distance $y$ and the system is made to perform SHM in vertical plane
Here, the springs have different force constants. Moreover the increase in their length is same. Therefore, the load is distributed equally between the springs. Hence, the restoring force developed in each spring is different.
If $F_{1}$ and $F_{2}$ are the restoring forces set up due to extension of springs, then
$F_{1}=-k_{1} y$ and $F_{2}=k_{2} y$
Also the total restoring force
$F=F_{1}+F_{2}$
$F=F_{1}+F_{2}$
$F=-k_{1} y-k_{2} y$
$F=-\left(k_{1}+k_{2}\right) y$
If $k^{\prime}$ is the is equivalent spring constant then
$F=-k^{\prime} y$ thus
$k^{\prime}=k_{1}+k_{2}$
Frequency and periodic time will be same as given in (i)
Q) A small mass $m$ is fastened to a vertical wire, which is under tension $T$. What will be the
 natural frequency of vibration of the mass if it is displaced laterally a slight distance and then released?

Solution


Sphere is displaced by a very small distance x thus angle formed is also small
Using trigonometry we can find Restoring force $\mathrm{F}_{b}=\mathrm{T} \sin \theta=T(x / b)$
Now restoring force per unit displacement $\mathrm{k}_{1}=\mathrm{F}_{\mathrm{b}} / \mathrm{x}=\mathrm{T} / \mathrm{b}$ $\mathrm{F}_{\mathrm{c}}=\mathrm{T} \sin \theta=\mathrm{T}(\mathrm{x} / \mathrm{c})$
Now restoring force per unit displacement $\mathrm{k}_{2}=\mathrm{F}_{\mathrm{c}} / \mathrm{x}=\mathrm{T} / \mathrm{c}$
Since both restoring forces are parallel combination

$$
\begin{gathered}
f=\frac{1}{2 \pi} \sqrt{\frac{\frac{T}{b}+\frac{T}{c}}{m}} \\
f=\frac{1}{2 \pi} \sqrt{\frac{T}{m}\left(\frac{1}{b}+\frac{1}{c}\right)} \\
f=\frac{1}{2 \pi} \sqrt{\frac{T}{m}\left(\frac{b+c}{b c}\right)}
\end{gathered}
$$

Q) The spring has a force constant $k$. the pulley is light and smooth while the spring and wew the string are light. If the block of mass ' $m$ ' is slightly displaced vertically and released, find the period of vertical oscillation

Solution: When mass $m$ is pulled by distance of $x$, increase in length of spring is $x / 2$ as explained below


Before pulling mass m spring was in equilibrium
$2 T_{0}=k x_{0}$
And $\mathrm{T}_{0}=\mathrm{mg}$
Thus $2 \mathrm{mg}=\mathrm{kx} \mathrm{x}_{0}$
When spring stretch by $\mathrm{x} / 2$ then tension in spring is T

$$
\begin{gathered}
F=k\left(x_{0}+\frac{x}{2}\right) \\
\text { Or } 2 T=k\left(x_{0}+\frac{x}{2}\right) \\
2 T=k x_{0}+\frac{k x}{2} \\
2 T=2 T_{0}+\frac{k x}{2} \\
2 T-2 T_{0}=\frac{k x}{2} \\
T-T_{0}=\frac{k x}{4}
\end{gathered}
$$

Restoring force on mass m is $\mathrm{T}-\mathrm{T}_{0}$ which is proportional to displacement

Thus restoring force constant $\mathrm{k}^{\prime}=\mathrm{k} / 4$

Period of oscillation

$$
T=2 \pi \sqrt{\frac{m}{k^{\prime}}}=2 \pi \sqrt{\frac{m}{\frac{k}{4}}}=4 \pi \sqrt{\frac{m}{k}}
$$

## Differential equation of simple harmonic motion

According to Newton's second law of motion

$$
F=m a=m \frac{d^{2} y(t)}{d t^{2}}
$$

Comparing this with $\mathrm{F}=-\mathrm{ky}(\mathrm{t})$

$$
\begin{aligned}
& m \frac{d^{2} y(t)}{d t^{2}}=-k y(t) \\
& \frac{d^{2} y(t)}{d t^{2}}=-\frac{k}{m} y(t) \\
& \frac{d^{2} y(t)}{d t^{2}}=-\omega^{2} y(t) \\
& \frac{d^{2} y(t)}{d t^{2}}+\omega^{2} y(t)=0
\end{aligned}
$$

This is the second order differential equation of the simple harmonic motion. The solution of this equation is of the type
$Y(t)=A \sin \omega t$ or $y(t)=B \cos \omega t$
Or any linear combination of sine and cosine function
$Y(t)=A \sin \omega t+B \cos \omega t$

## Solved Numerical

Q) The SHM is represented by $y=3 \sin 314 t+4 \cos 314 t$. $y$ in cm and in $t$ in second. Find the amplitude, epoch the periodic time and the maximum velocity of SHO
Solution
$Y=A \sin (\omega t+\varphi)$
$Y=A \cos \varphi \sin \omega t+A \sin \varphi \cos \omega t$
Comparing with $3 \sin 314 t+4 \cos 314 t$
$3=A \cos \varphi$ and $4=A \sin \varphi$
$\therefore \mathrm{A}^{2} \cos ^{2} \varphi+\mathrm{A}^{2} \sin ^{2} \varphi=3^{2}+4^{2}$
$A^{2}=25$
$A=5 \mathrm{~cm}$
The initial phase (epoch) is obtained as

$$
\begin{gathered}
\tan \varphi=\frac{\sin \varphi}{\cos \varphi}=\frac{4}{3} \\
\varphi=\tan ^{-1}\left(\frac{4}{3}\right) \\
\varphi=53^{\circ} 8^{\prime}
\end{gathered}
$$

Now

$$
T=\frac{2 \pi}{\omega \pi}=\frac{2}{314}=0.02 \mathrm{~s}
$$

Maximum velocity
$\mathrm{V}_{\max }=\omega \mathrm{A}=314 \times 5=1570 \mathrm{~cm} / \mathrm{s}$
Q) The vertical motion of a ship at sea is described by the equation $\frac{d^{2} x}{d t^{2}}=-4 x$ where x in metre is the vertical height of the ship above its mean position. If it oscillates through a total distance of 1 m in half oscillation, find the greatest vertical speed and the greatest vertical acceleration.
Solution
Comparing given with standard equation for oscillation we get

$$
\omega^{2}=4 \text { or } \omega=2
$$

given amplitude $A=1 / 2$
$V_{\text {max }}=\omega A=2 \times(1 / 2)=1 \mathrm{~m} / \mathrm{s}$
$A_{\max }=\omega^{2} A=2^{2} \times(1 / 2)=2 \mathrm{~m} / \mathrm{s}^{2}$

## Total Mechanical Energy in Simple Harmonic Oscillator

The total energy $(E)$ of an oscillating particle is equal to the sum of its kinetic energy and potential energy if conservative force acts on it.

## Kinetic energy

Kinetic energy of the particle of mass $m$ is

$$
K=\frac{1}{2} m v^{2}=\frac{1}{2} m \omega^{2}\left(A^{2}-x^{2}\right)---e q(1)
$$

## Potential energy

From definition of SHM $F=-k x$ the work done by the force during the small displacement $d x$ is $d W=-F . d x=-(-k x) d x=k x d x$
$\therefore$ Total work done for the displacement $x$ is,

$$
W=\int d w=\int_{0}^{x} k x d x
$$

$k=\omega^{2} m$

$$
W=\int_{0}^{x} m \omega^{2} x d x=\frac{1}{2} m \omega^{2} x^{2}
$$

This work done is stored in the body as potential energy

$$
U=\frac{1}{2} m \omega^{2} x^{2}---(2)
$$

Total energy $E=K+U$

$$
\begin{gathered}
E=\frac{1}{2} m \omega^{2}\left(A^{2}-x^{2}\right)+\frac{1}{2} m \omega^{2} x^{2} \\
E=\frac{1}{2} m \omega^{2} A^{2}
\end{gathered}
$$

## Special cases

(i) When the particle is at the mean position $x=0$, from eqn (1) it is known that kinetic energy is maximum total energy is wholly kinetic $K_{\max }=\frac{1}{2} m \omega^{2} A^{2}$ and from eqn. (2) it is known that potential energy is zero.
(ii) When the particle is at the extreme position $\mathrm{y}=+a$, from eqn. (1) it is known that kinetic energy is zero and from eqn. (2) it is known that Potential energy is maximum. Hence the total energy is wholly potential. $U_{\max }=\frac{1}{2} m \omega^{2} A^{2}$
(iii)When $\mathrm{y}=\mathrm{A} / 2$

$$
\begin{gathered}
K=\frac{1}{2} m \omega^{2}\left[A^{2}-\left(\frac{A}{2}\right)^{2}\right] \\
K=\frac{3}{4}\left(\frac{1}{2} m \omega^{2} A^{2}\right) \\
K=\frac{3}{4}(E) \\
U=\frac{1}{2} m \omega^{2}\left(\frac{A}{2}\right)^{2} \\
U=\frac{1}{4}\left(\frac{1}{2} m \omega^{2} A^{2}\right)
\end{gathered}
$$

If the displacement is half of the amplitude $K$ and $U$ are in the ratio $3: 1$,

## Solved numerical

Q) If a particle of mass 0.2 kg executes SHM of amplitude 2 cm and period of 6 sec find(iOthe total mechanical energy at any instant (ii) kinetic energy and potential energies when the displacement is 1 cm
Solution:

$$
\omega=\frac{2 \pi}{T}=\frac{2 \pi}{6}
$$

(i)Total mechanical energy at any instant is given by

$$
\begin{gathered}
E=\frac{1}{2} m \omega^{2} A^{2} \\
E=\frac{1}{2}(0.2)\left(\frac{2 \pi}{6}\right)^{2}\left(2 \times 10^{-2}\right)^{2} \\
E=0.1 \times \frac{4 \pi^{2}}{36} \times 4 \times 10^{-4} \\
E=4.39 \times 10^{-5} \mathrm{~J}
\end{gathered}
$$

(ii)K.E. at the instant when displacement x is given by

$$
\begin{gathered}
K=\frac{1}{2} m \omega^{2}\left(A^{2}-x^{2}\right) \\
K=\frac{1}{2}(0.2)\left(\frac{2 \pi}{6}\right)^{2}\left(4 \times 10^{-4}-1 \times 10^{-4}\right) J \\
K=3.29 \times 10^{-5} \mathrm{~J}
\end{gathered}
$$

P.E. energy at that instant $=$ Total energy - K.E
$=(4.39-3.29) \times 10^{-5}$
$=1.1 \times 10^{-5} \mathrm{~J}$

## Angular Simple Harmonic Motion

A body to rotate about a given axis can make angular oscillations. For example, a wooden stick nailed to a wall can oscillate about its mean position in the vertical plane The conditions for an angular oscillation to be angular harmonic motion are
(i)When a body is displaced through an angle from the mean position, the resultant torque is proportional to the angle displaced
(ii)This torque is restoring in nature and it tries to bring the body towards the mean position


If the angular displacement of the body at an instant is $\theta$, then resultant torque on the body
$\tau=-k \theta$
if the momentum of inertia is $I$, the angular acceleration is

$$
\begin{gathered}
\alpha=\frac{\tau}{I}=-\frac{k}{I} \theta \\
\text { Or } \\
\frac{d^{2} \theta}{d t^{2}}=-\omega^{2} \theta \quad---e q(1)
\end{gathered}
$$

Here

$$
\omega=\sqrt{\frac{k}{I}}
$$

Solution of equation (1) is

$$
\theta=\theta_{0} \sin (\omega t+\varphi)
$$

Where $\theta_{0}$ is the maximum angular displacement on either side. Angular velocity at time ' t ' is given by

$$
\omega=\frac{d \theta}{d t}=\theta_{0} \omega \cos (\omega t+\varphi)
$$

## Physical pendulum

An rigid body suspended from a fixed support constitutes a physical pendulum.
As shown in figure is a physical pendulum. A rigid body
 is suspended through a hole at O . When the centre of mass $C$ is vertically below $O$ at a distance of ' $I$ ', the body may remain at rest.
The body is rotated through an angle $\theta$ about a horizontal axis OA passing through $O$ and perpendicular to the plane of motion The torque of the forces acting on the body, about the axis $O A$ is $\tau=m g l \sin \theta$, here $l=$ OC
If momentum of inertia of the body about $O A$ is $I$, the angular acceleration becomes

$$
\alpha=\frac{\tau}{I}=-\frac{m g l}{I} \sin \theta
$$

For small angular displacement $\sin \theta=\theta$

$$
\alpha=-\left(\frac{m g l}{I}\right) \theta
$$

Comparing with $\alpha=-\omega^{2} \theta$

$$
T=\frac{2 \pi}{\omega}=2 \pi \sqrt{\frac{I}{m g l}}
$$

## Solved Numerical

Q) A uniform meter stick is suspended through a small-hole at the 10 cm mark. Find the time period of small oscillations about the point of suspension
Solution:

Let the mass of stick be $M$. The moment of inertia of the stick about the axis of rotation through the point of suspension is

$$
I=\frac{M l^{2}}{12}+M d^{2}=M\left(\frac{l^{2}}{12}+d^{2}\right)
$$

Centre of mass of the stick is at 50 cm from top thus it is at distance 40 cm from the small hole thus $d=40 \mathrm{~cm}$, given length of stick $=1 \mathrm{~m}$
Time period

$$
\begin{array}{r}
T=2 \pi \sqrt{\frac{I}{M g d}} \\
T=2 \pi \sqrt{\frac{M\left(\frac{l^{2}}{12}+d^{2}\right)}{M g d}} \\
T=2 \pi \sqrt{\frac{\left(\frac{l^{2}}{12}+d^{2}\right)}{g d}} \\
T=2 \times 3.14 \sqrt{\frac{1}{\frac{12}{9}+(0.4)^{2}}} \\
\quad \mathrm{~T}=1.56 \mathrm{sec}
\end{array}
$$

## Torsional pendulum

In torsional pendulum, an extended body is suspended by a light thread or wire. The body is rotated through an angle about the wire as the axis of rotation.
The wire remains vertical during this motion but a twist ' $\theta$ ' is produced in the wire. The twisted wire exerts a restoring torque on the body, which is proportional to the angle of the twist.
$\tau \propto-\theta ; \tau=-k \theta ; k$ is proportionality constant and is called torsional constant of the wire. If 1 be the moment of inertia of the body about vertical axis, the angular acceleration is

$$
\begin{gathered}
\alpha=\frac{\tau}{I}=\frac{-k}{I} \theta=-\omega^{2} \theta \\
\omega=\sqrt{\frac{k}{I}}
\end{gathered}
$$

Time period $T=\frac{2 \pi}{\omega}=2 \pi \sqrt{\frac{I}{k}}$

## PHYSICS NOTES

Q) The moment of inertia of the disc used in a torsional pendulum about the suspension wire is $0.2 \mathrm{~kg}-\mathrm{m}^{2}$. It oscillates with a period of 2 s . Another disc is placed over the first one and the time period of the system becomes 2.5 s . Find the moment of inertia of the second disc about the wire

## Solution

Let the torsional constant of the wire be k

$$
\begin{gathered}
T=2 \pi \sqrt{\frac{I}{k}} \\
2=2 \pi \sqrt{\frac{0.2}{k}}---e q(1)
\end{gathered}
$$

When a second disc having moment of inertia $I_{1}$ about the wire is added, the time period is

$$
2.5=2 \pi \sqrt{\frac{0.2+I_{1}}{k}}---e q(2)
$$

From eq(1) and (2)

$$
\mathrm{I}_{1}=0.11 \mathrm{~kg}-\mathrm{m}^{2}
$$

## Two Body System

In a two body oscillations, such as shown in the figure, a spring connects two objects, each of which is free to move. When the objects are displaced and released, they both oscillate. The relative separation $x_{1}-x_{2}$ gives the length of the spring at any time. Suppose its unscratched length is $L$; then $x=\left(x_{1}-x_{2}\right)-L$ is the change in length of the spring, and $F=k x$ is the magnitude of the force exerted on each particle by the spring as shown in figure.
Applying Newton's second law separately to the two particles, taking force component along the $x$-axis, we get

$$
m_{1} \frac{d^{2} x_{1}}{d t^{2}}=-k x \text { and } m_{2} \frac{d^{2} x_{2}}{d t^{2}}=+k x
$$

Multiplying the first of these equations by $m_{2}$ and the second by $m_{1}$ and then subtracting,

$$
m_{1} m_{2} \frac{d^{2} x_{1}}{d t^{2}}-m_{1} m_{2} \frac{d^{2} x_{2}}{d t^{2}}=-m_{2} k x-m_{1} k x
$$

This can be written as

$$
\begin{gathered}
\frac{m_{1} m_{2}}{\left(m_{1}+m_{2}\right)} \frac{d^{2}}{d t^{2}}\left(x_{1}-x_{2}\right)=-k x \\
\mu \frac{d^{2}}{d t^{2}}\left(x_{1}-x_{2}\right)=-k x---e q(1)
\end{gathered}
$$

Here $\mu$ is known as reduced mass and has dimension of mass

$$
\mu=\frac{m_{1} m_{2}}{\left(m_{1}+m_{2}\right)}
$$

## Since L is constant

$$
\frac{d^{2}}{d t^{2}}\left(x_{1}-x_{2}\right)=\frac{d^{2}}{d t^{2}}(x+L)=\frac{d x}{d t}
$$

Equation (i) becomes

$$
\begin{aligned}
\mu \frac{d^{2} x}{d t^{2}} & =-k x \\
\frac{d^{2} x}{d t^{2}} & =-\frac{k}{\mu} x
\end{aligned}
$$

Thus periodic time

$$
T=2 \pi \sqrt{\frac{\mu}{k}}
$$

Q) Two balls with masses $m_{1}=1 \mathrm{~kg}$ and $m_{2}=2 \mathrm{kgkg}$ are slipped on a thin smooth horizontal
 rod. The balls are interconnected by a light spring of spring constant $24 \mathrm{~N} / \mathrm{m}$. the left hand ball is imparted the initial velocity $\mathrm{v}_{1}=12 \mathrm{~cm} / \mathrm{s}$.
Find (a) the oscillation frequency of the system (b) the energy and amplitude of oscillation

## Solution

Reduced mass

$$
\mu=\frac{m_{1} m_{2}}{\left(m_{1}+m_{2}\right)}=\frac{1 \times 2}{1+2}=\frac{2}{3} \mathrm{~kg}
$$

Frequency
$f=1 / T$

$$
\begin{gathered}
f=\frac{1}{2 \pi} \sqrt{\frac{k}{\mu}} \\
f=\frac{1}{2 \times 3.14} \sqrt{\frac{24 \times 3}{2}}=0.955 \mathrm{sec}
\end{gathered}
$$

Initial velocity given to mass $m_{1}$ is $v_{1}$
For undamped oscillation, this initial energy remains constant
Hence total energy of S.H.M. of two balls is given as

$$
E=\frac{1}{2} \mu v_{1}^{2}
$$

If amplitude of oscillation is A then

$$
\frac{1}{2} \mu v_{1}^{2}=\frac{1}{2} k A^{2}
$$

$$
\begin{gathered}
A=\sqrt{\frac{\mu}{k}} v_{1} \\
A=\sqrt{\frac{2}{3 \times 24}} \times 0.12=0.02 \mathrm{~m}
\end{gathered}
$$

Or $A=2 \mathrm{~cm}$

## Damped oscillation

Experimental studies showed that the resistive force acting on the oscillator in a fluid medium depends upon the velocity of the oscillator
Thus resistive force or damping force acting on the oscillator is
$F_{d} \propto v$
$\therefore \mathrm{F}_{\mathrm{d}}=-\mathrm{nv}$
Here $b$ is damping constant and has SI units $\mathrm{kg} /$ second. The negative sign indicates that the force $F_{d}$ opposes the motion
Thus, a damped oscillator oscillate under the influence of the following forces
(i)Restoring force $F_{x}=-k x$ and
(ii) Resistive force $F_{d}=-b v$

Net force $F=F_{x}+F_{d}$
According to second law of motion
$m a=-k y-b v$

$$
\begin{array}{r}
m \frac{d^{2} x}{d t^{2}}=-\mathrm{ky}-\mathrm{b} \frac{\mathrm{dx}}{\mathrm{dt}} \\
\frac{d^{2} x}{d t^{2}}+\frac{\mathrm{b}}{\mathrm{~m}} \frac{\mathrm{dx}}{\mathrm{dt}}+\frac{\mathrm{k}}{\mathrm{~m}} \mathrm{y}=0 \quad-- \tag{1}
\end{array}
$$

Solution of this equation is

$$
x(t)=A e^{-b t / 2 m} \sin \left(\omega^{\prime} t+\varphi\right)
$$

Here $A e^{-b t / 2 m}$ is the amplitude of the damped oscillation and decreases exponentially with time
The angular frequency $\omega^{\prime}$ of the damped oscillator is given by

$$
\omega^{\prime}=\sqrt{\frac{k}{m}-\frac{b^{2}}{4 m^{2}}}
$$

Energy of oscillator

$$
E=\frac{1}{2} k A^{2} e^{-b t / m}
$$

Above equation is valid only if $b \ll \sqrt{ }(\mathrm{~km})$

## PHYSICS NOTES

## Natural Oscillations

When a system capable of oscillating is given some initial displacement from its equilibrium position and left free (i.e. in absence of any external force) it begins to oscillate. Thus the oscillations performed by it in absence of any resistive forces are known as natural oscillations. The frequency of natural oscillations is known as natural frequency $f_{0}$ and corresponding angular frequency is denoted by $\omega^{0}$

## Forced Oscillation

Oscillations of the system under the influence of an external periodic force are forced oscillation
Consider an external periodic force $F=F_{0} \sin \omega t$ acting on the system which is capable to oscillate
Equation for oscillation can be witten as

$$
\begin{aligned}
& \qquad m \frac{d^{2} x}{d t^{2}}+\mathrm{b} \frac{\mathrm{dx}}{\mathrm{dt}}+\mathrm{ky}=\mathrm{F}_{0} \sin \omega \mathrm{t} \\
& \qquad \frac{d^{2} x}{d t^{2}}+\frac{\mathrm{b}}{\mathrm{~m}} \frac{\mathrm{dx}}{\mathrm{dt}}+\frac{\mathrm{k}}{\mathrm{~m}} \mathrm{y}=\frac{\mathrm{F}_{0}}{\mathrm{~m}} \sin \omega \mathrm{t} \\
& \text { The solution of equation is given by } \\
& \qquad \mathrm{X}=\mathrm{A} \sin (\omega \mathrm{t}+\varphi)
\end{aligned}
$$

Here, $A$ and $\varphi$ are the constants of the solution they are found as,

$$
\begin{gathered}
A=\frac{F_{0}}{\left[m^{2}\left(\omega_{0}^{2}-\omega^{2}\right)^{2}+b^{2} \omega^{2}\right]^{1 / 2}} \\
\varphi=\tan ^{-1} \frac{\omega x}{v_{0}}
\end{gathered}
$$

Here $m$ is the mass of oscillator, $v_{0}$ is velocity of oscillator, $x$ is the displacement of oscillator.
(i) for small damping factor

$$
m\left(\omega_{0}^{2}-\omega^{2}\right) \gg b \omega
$$

Equation for amplitude becomes

$$
A=\frac{F_{0}}{m^{2}\left(\omega_{0}^{2}-\omega^{2}\right)^{2}}
$$

(ii) For large damping factor

$$
\begin{gathered}
b \omega \gg m\left(\omega_{0}^{2}-\omega^{2}\right) \\
A=\frac{F_{0}}{b \omega}
\end{gathered}
$$

If when value of $\omega$ approaches $\omega_{0}$ the amplitude becomes maximum. This phenomenon is known as resonance. The value of $\omega$ for which resonance occurs is known as the resonant frequency.

## PHYSICS NOTES

Solved numerical
Calculate the time during which the amplitude becomes $A / 2^{n}$ in case of damped oscillations, where $A=$ initial amplitude
Solution:

$$
A(t)=A e^{-b t / 2 m}
$$

But $A(t)=A / 2^{n}$

$$
\therefore \frac{A}{2^{n}}=A e^{-b t / 2 m}
$$

Taking log to the base e on both sides

$$
\begin{aligned}
& \frac{b t}{2 m}=n \ln 2 \\
& t=\frac{2 m n \ln 2}{b t}=\frac{2 m n}{b t}(0.693)
\end{aligned}
$$

To become a member of www.gneet.com give call on 09737240300

